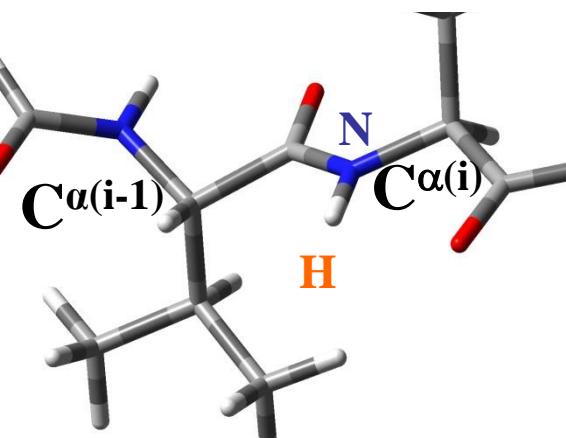
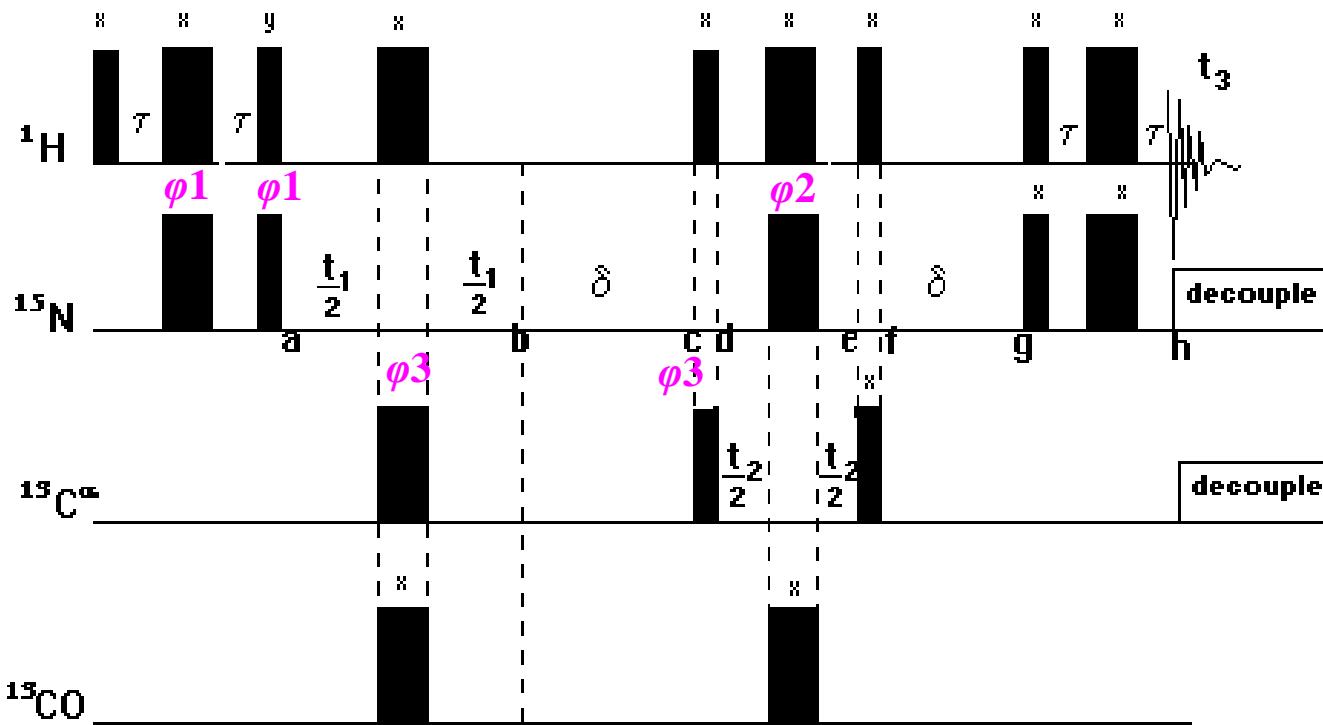


The 3D-HNCA on a ^{15}N - ^{13}C labeled sample:



$$\varphi 1 = x, -x$$

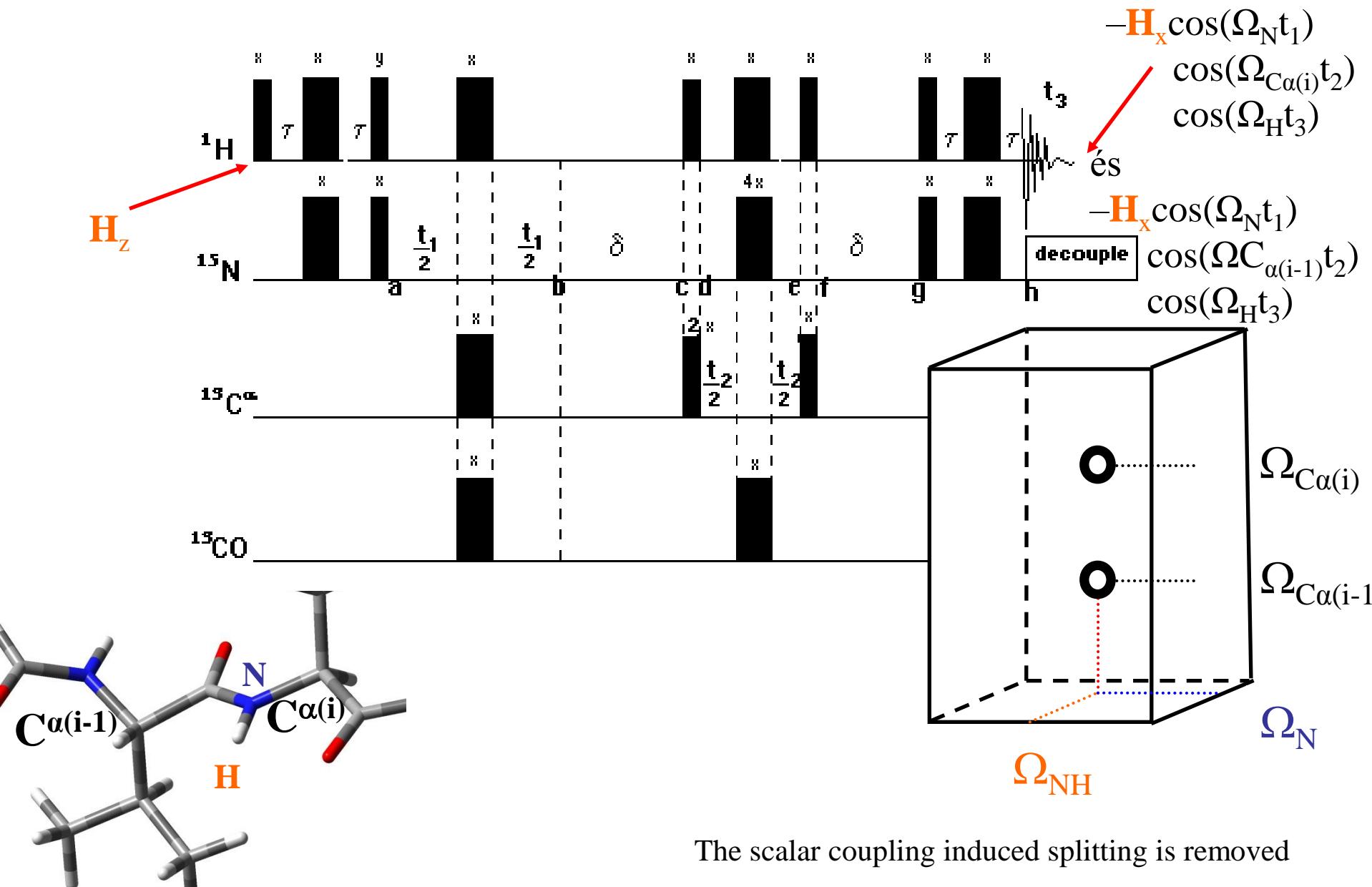
$$\varphi 2 = 4(x), 4(y), 4(-x), 4(-y)$$

$$\varphi 3 = 2(x), 2(-x)$$

$$receiver \quad x = x, -x, -x, x, -x, x, x, -x$$

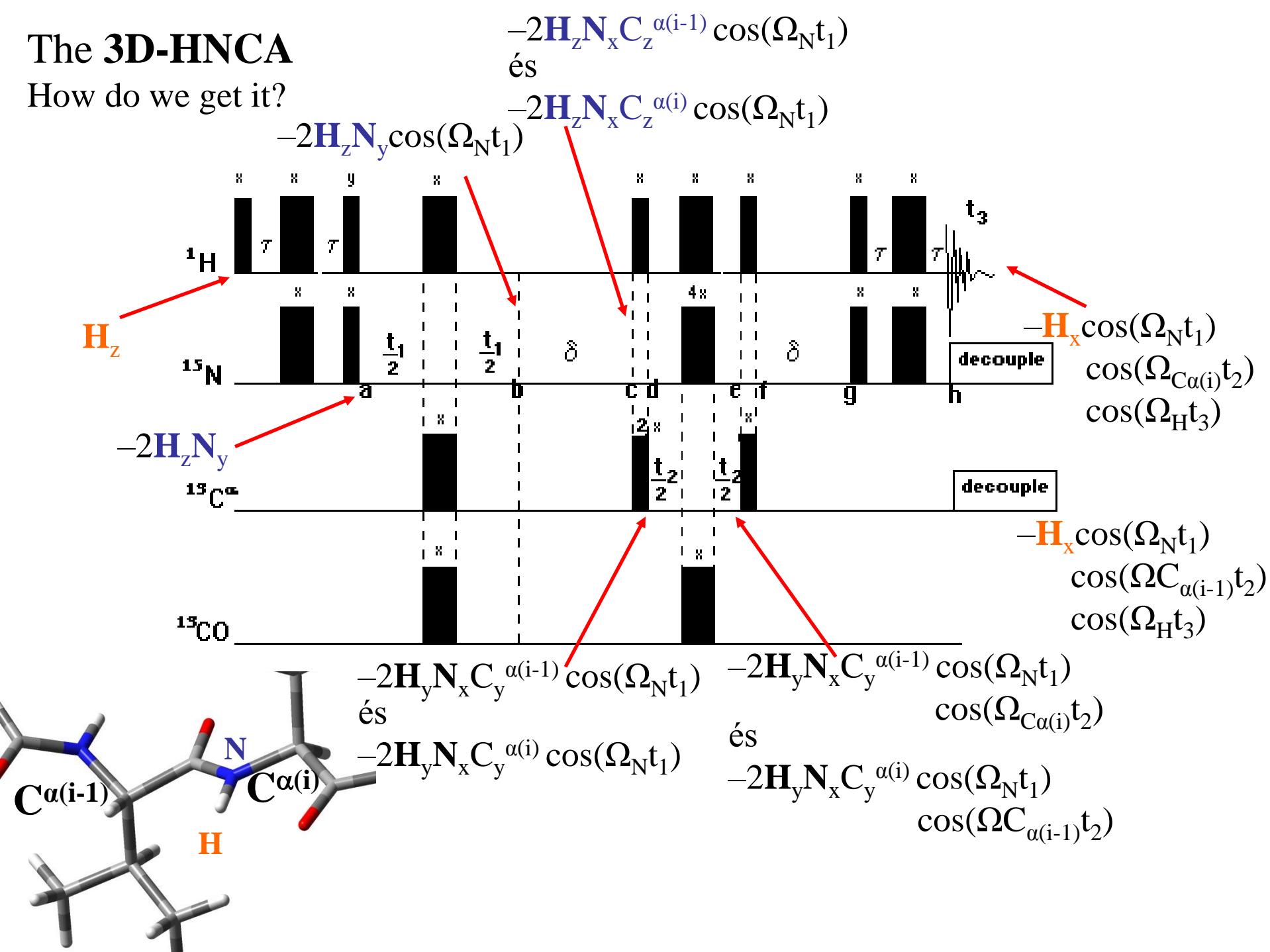
The 3D-HNCA on a ^{15}N , ^{13}C labeled sample:

What do we get at the end?



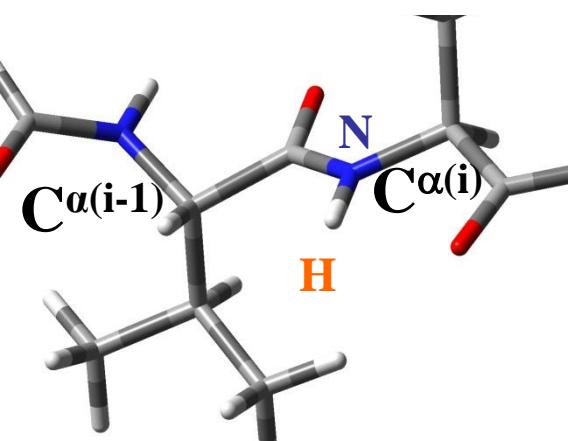
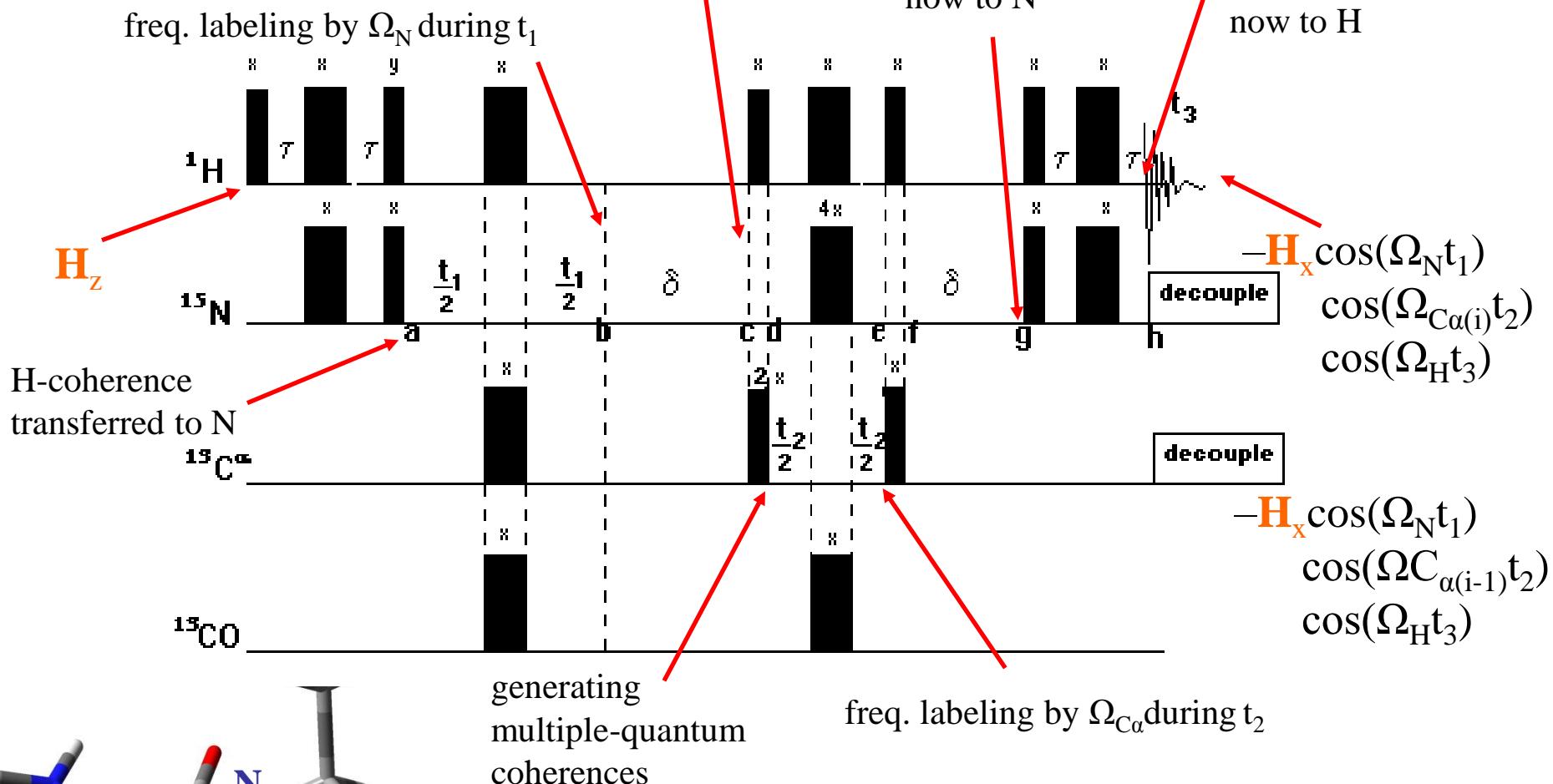
The 3D-HNCA

How do we get it?



The 3D-HNCA

How do we get it?



A. module (getting started)

$\sigma[0]$ "The first INEPT module"

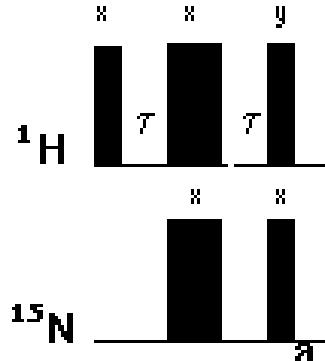
$$\hat{H} = H_x(\pi/2)$$

$$\hat{H} = \text{echo (homo)}$$



$$+H_z$$

$$-H_y$$



$$\hat{H} = H_y(\pi/2)$$

$$\hat{H} = N_x(\pi/2)$$

$\sigma[a]$ "end of this INEPT"

$$-2H_xN_z$$



$$+2H_zN_z$$



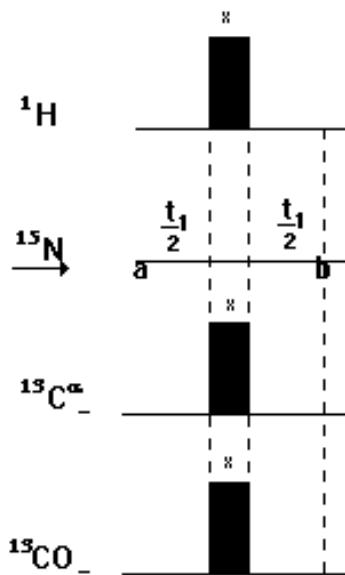
$$-2H_zN_y$$

B. module (frequency labeling by N)

$\sigma[a] \rightarrow \sigma[b]$ "frequency labeling" by N^{15}

The echo module decouples N from all the other nuclei (e.g. H, C^α , C').

In other words, all scalar couplings between ^{15}N and $^1H^{NH}$, ^{13}Ca and ^{13}CO spins are removed by the 180° refocusing pulses (echo) positioned at the middle of t_1 .



$\sigma[a]$ "at the beginning of t_1 "

\hat{H} = echo (hetero)

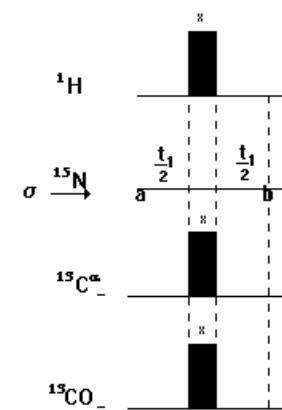
$\sigma[b]$ "at the end of t_1 "

$$-2\mathbf{H}_z \mathbf{N}_y$$



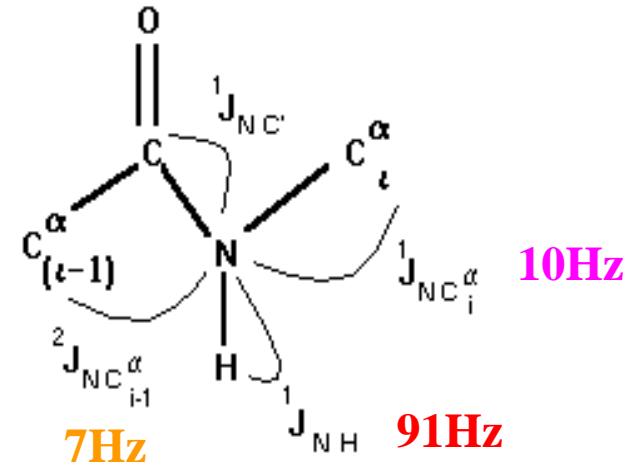
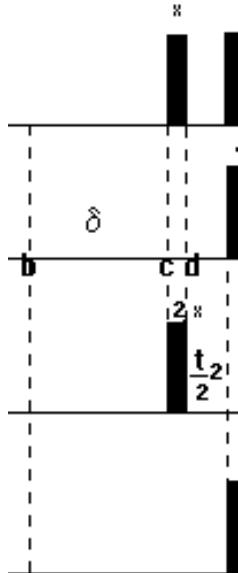
$$+2\mathbf{H}_z \mathbf{N}_y \cos(\Omega_N t_1) - 2\mathbf{H}_z \mathbf{N}_x \sin(\Omega_N t_1)$$

memo : $-2\mathbf{H}_z \mathbf{N}_x \sin(\Omega_N t_1)$ is phase cycled out, as it would evolve into a $+4\mathbf{H}_z \mathbf{N}_y \mathbf{C}_z$ type anti-phase „tri-spin” magnetization.



C. module (transferring coherence from N to Cα(s))

$\sigma[b]$ from the beginning of delay δ term $+2\mathbf{H}_z \mathbf{N}_y \cos(\Omega_N t_1)$ is to be considered only *memo*: during the delay-time δ "all" couplings ($^1J_{HN}$, $^1J_{NC}$, $^1J_{NC}^\alpha$, $^2J_{NC}^\alpha$) are active.



$$\hat{H} = 2H_z N_z ({}^1J_{NH} \pi \delta)$$

$$+ 2H_z N_y \cos(\Omega_N t_1)$$

$$+ 2H_z N_y \cos(\Omega_N t_1) \cos({}^1J_{NH} \pi \delta) + 2H_z N_x \cos(\Omega_N t_1) \sin({}^1J_{NH} \pi \delta)$$



memo 3: If δ is set to a value $k/{}^1J_{HN}$ (where $k = 1, 2, \dots$) [$1/91, 2/91, \text{etc.}$] then ${}^1J_{NH} \pi \delta \rightarrow k^* \pi$ [$1\pi, 2\pi, \text{etc.}$]. As $\sin(k^* \pi) = 0$, the $+ 2H_z N_x \cos(\Omega_N t_1) \sin({}^1J_{NH} \pi \delta)$ term must vanish.

Thus, $+ 2H_z N_y \cos(\Omega_N t_1) \cos({}^1J_{NH} \pi \delta)$ is to be considered only:

$$\hat{H} = 2N_z C_z ({}^1J_{NC} \alpha^{(i)} \pi \delta)$$

$$\begin{aligned} \hat{H} = 2N_z C_z ({}^2J_{NC} \alpha^{(i-1)} \pi \delta) \\ + 2H_z N_y \end{aligned}$$

$$\begin{aligned} + 2H_z N_y \\ - 4H_z N_x C \alpha^{(i-1)} z \\ - 4H_z N_x C \alpha^{(i)} z \\ - 8H_z N_y C \alpha^{(i)} z C \alpha^{(i-1)} z \end{aligned}$$

$$+ 2H_z N_y$$



$$- 4H_z N_x C \alpha^{(i)} z$$



$$- 4H_z N_x C \alpha^{(i)} z$$



$$- 8H_z N_y C \alpha^{(i)} z C \alpha^{(i-1)} z$$

$$\begin{aligned} * \cos({}^1J_{NC} \alpha^{(i)} \pi \delta) \\ * \cos({}^1J_{NC} \alpha^{(i)} \pi \delta) \\ * \sin({}^1J_{NC} \alpha^{(i)} \pi \delta) \\ * \sin({}^1J_{NC} \alpha^{(i)} \pi \delta) \end{aligned}$$

$$\begin{aligned} * \cos({}^2J_{NC} \alpha^{(i-1)} \pi \delta) \\ * \sin({}^2J_{NC} \alpha^{(i-1)} \pi \delta) \\ * \cos({}^2J_{NC} \alpha^{(i-1)} \pi \delta) \\ * \sin({}^2J_{NC} \alpha^{(i-1)} \pi \delta) \end{aligned}$$

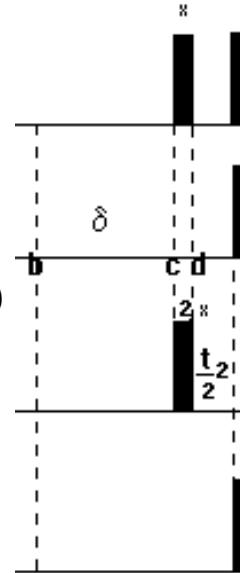
a 2-spin term

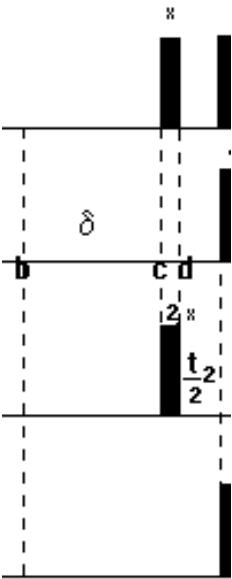
a 3-spin term

a 3-spin term

a 4-spin term

memo: recall that ${}^1J_{NC} \alpha^{(i)} \neq {}^2J_{NC} \alpha^{(i-1)}$ as $10 \neq 7 \text{Hz}$





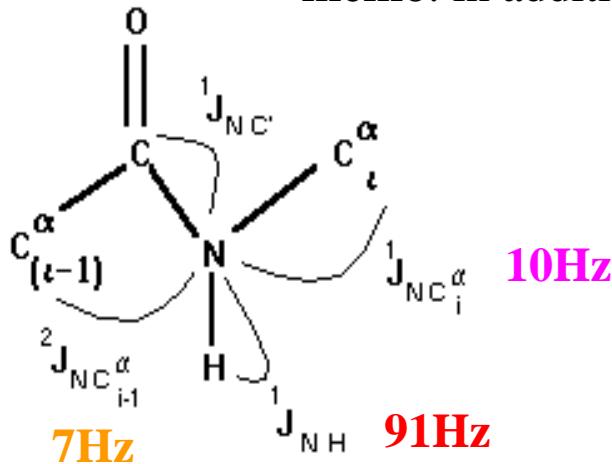
question: how to set optimally delay time δ ?

answer: the $\cos(^1J_{NC}\alpha^{(i)}\pi\delta)$ * $\sin(^2J_{NC}\alpha^{(i-1)}\pi\delta)$ and the $\sin(^1J_{NC}\alpha^{(i)}\pi\delta)$ * $\cos(^2J_{NC}\alpha^{(i-1)}\pi\delta)$ terms are called as the "transfer" functions; terms to be maximized.

Considering that $^1J_{NC}\alpha^{(i)}$ and $^2J_{NC}\alpha^{(i-1)}$ are 10 and 7 Hz, respectively (with $^1J_{NH} = 91$ Hz) both transfer functions have a maximum at $\delta \sim 22$ ms.

So one can now **set δ at a sensible value (~ 22 ms)** for globular proteins.

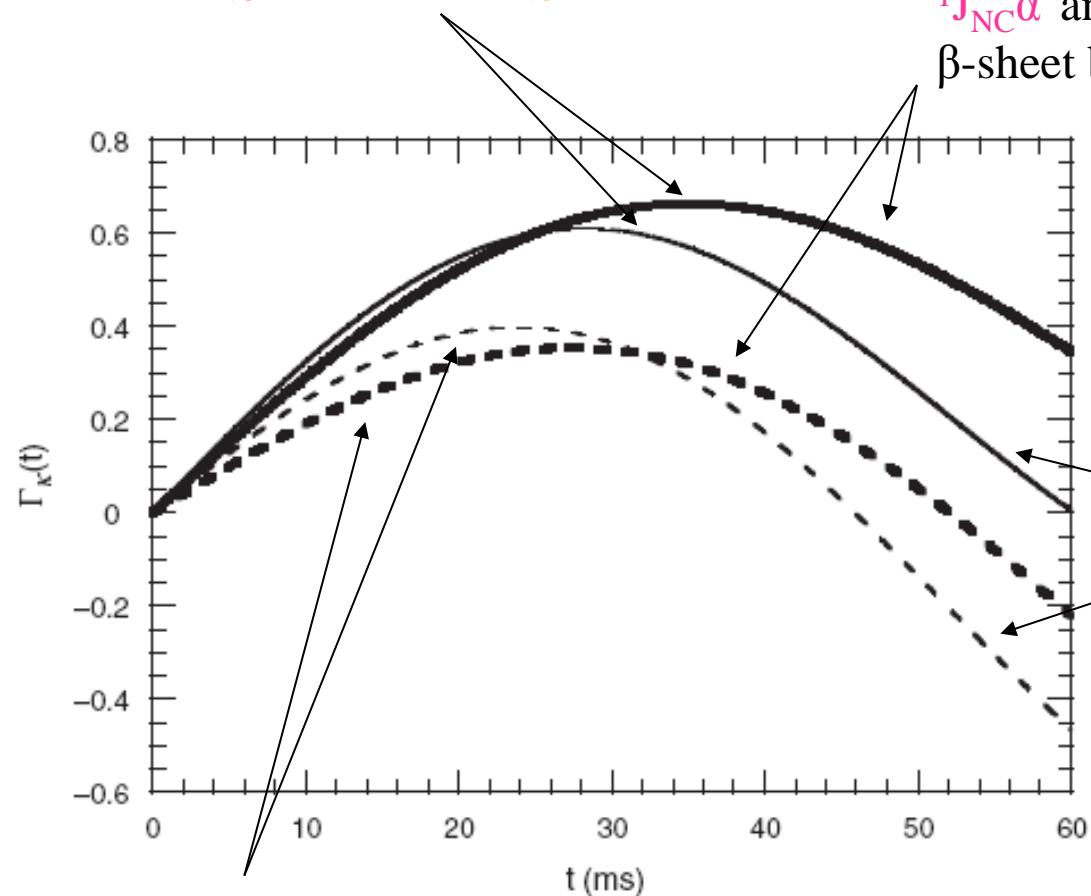
memo: in addition, if $\delta \sim 22$ ms, then $\cos(^1J_{NH}\pi\delta) \sim \cos(2,002\pi) \sim 1$.



$^1J_{NC}$ refocused with the 180° suite centered in the middle of t_2 (or more typically it is decoupled).

The two "transfer" functions, f_1 and f_2 , as function of the delay-time (δ)

$$f_1(\delta) = \cos(^1J_{NC}\alpha^{(i)}\pi\delta) \sin(^2J_{NC}\alpha^{(i-1)}\pi\delta)$$



$^1J_{NC}\alpha$ and $^2J_{NC}\alpha$ set for a
 β -sheet backbone structure

For α -helix $^1J_{NC}\alpha = 10.9\text{Hz}$
and $^2J_{NC}\alpha = 8.3\text{Hz}$

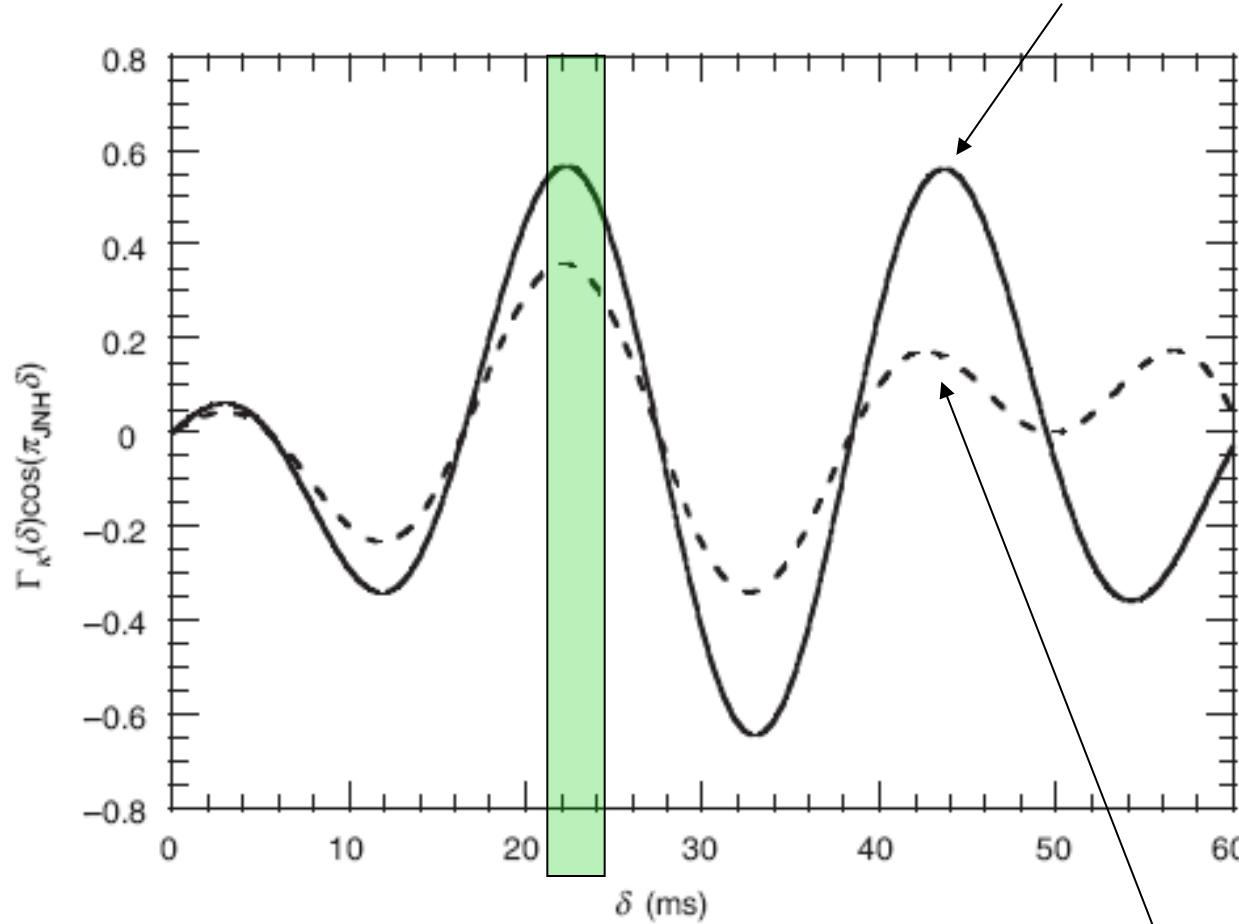
$^1J_{NC}\alpha$ and $^2J_{NC}\alpha$ set
for an α -helical
backbone structure

For α -helix $^1J_{NC}\alpha = 9.6\text{Hz}$
and $^2J_{NC}\alpha = 6.4\text{Hz}$

$$f_2(\delta) = \sin(^1J_{NC}\alpha^{(i)}\pi\delta) \cos(^2J_{NC}\alpha^{(i-1)}\pi\delta)$$

The two "transfer" functions, $f_1 \cos(^1J_{NH}\pi\delta)$ and $f_2 \cos(^1J_{NH}\pi\delta)$ as function of the delay-time (δ)

$$f_1(\delta) = \cos(^1J_{NC}\alpha^{(i)} \pi\delta) \sin(^2J_{NC}\alpha^{(i-1)} \pi\delta) \cos(^1J_{NH}\pi\delta)$$



One can set δ at a sensible value
(~ 22 ms) for globular proteins.

$$f_2(\delta) = \sin(^1J_{NC}\alpha^{(i)} \pi\delta) \cos(^2J_{NC}\alpha^{(i-1)} \pi\delta) \cos(^1J_{NH}\pi\delta)$$

memo 4: - since ${}^1J_{NC}\alpha^{(i)}$ and ${}^2J_{NC}\alpha^{(i-1)}$ couplings are different, they do not cancel out,
 - the anti-phase two spin ($+2H_zN_y$) and four-spin ($-8H_zN_yC\alpha(i)_zC\alpha(i-1)_z$) terms are removed by phase cycling during the next 90° on C (of phase $\varphi 3$).

Thus, the two 3-spin terms of interest are as follows:

$$\begin{array}{ll} -4H_zN_xC\alpha^{(i-1)}_z & * \cos({}^1J_{NC}\alpha^{(i)} \pi\delta) \\ -4H_zN_xC\alpha^{(i)}_z & * \sin({}^1J_{NC}\alpha^{(i)} \pi\delta) \end{array} \quad \begin{array}{ll} * \sin({}^2J_{NC}\alpha^{(i-1)} \pi\delta) \\ * \cos({}^2J_{NC}\alpha^{(i-1)} \pi\delta) \end{array} \quad \begin{array}{ll} * \cos({}^1J_{NH}\pi\delta) * \cos(\Omega_N t_1) \\ * \cos({}^1J_{NH}\pi\delta) * \cos(\Omega_N t_1) \end{array}$$

D. module (generating multiple-quantum coherences)

$$\sigma[c]$$

$$\hat{H} = H_x(\pi/2)$$

$$\hat{H} = C\alpha_x(\pi/2)$$

$$-4H_zN_xC_z$$



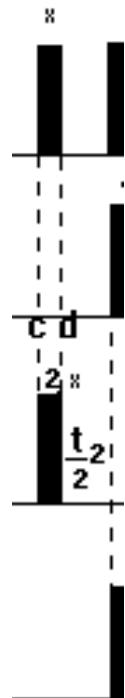
$$+4H_yN_xC_z$$



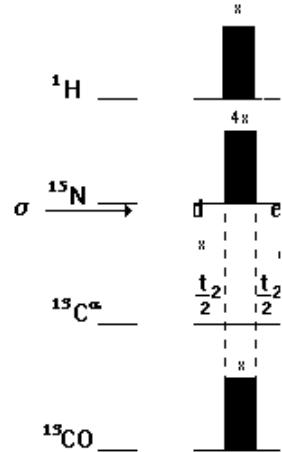
$$-4H_yN_xC_y$$

$$\sigma[d]$$

$$\begin{array}{ll} -4H_yN_xC\alpha^{(i-1)}_y * \cos({}^1J_{NC}\alpha^{(i)} \pi\delta) & * \sin({}^2J_{NC}\alpha^{(i-1)} \pi\delta) \\ -4H_yN_xC\alpha^{(i)}_y * \sin({}^1J_{NC}\alpha^{(i)} \pi\delta) & * \cos({}^2J_{NC}\alpha^{(i-1)} \pi\delta) \end{array} \quad \begin{array}{ll} * \cos({}^1J_{NH}\pi\delta) * \cos(\Omega_N t_1) \\ * \cos({}^1J_{NH}\pi\delta) * \cos(\Omega_N t_1) \end{array}$$



E. module (frequency labeling by C_α(s))



$\sigma[d]$ "at the beginning of t₂"

$$\begin{aligned} -4\mathbf{H}_y\mathbf{N}_x\mathbf{C}\alpha^{(i-1)}_y & * \cos(1J_{NC}\alpha^{(i)}\pi\delta) * \sin(2J_{NC}\alpha^{(i-1)}\pi\delta) * \cos(1J_{NH}\pi\delta) * \cos(\Omega_N t_1) \\ -4\mathbf{H}_y\mathbf{N}_x\mathbf{C}\alpha^{(i)}_y & * \sin(1J_{NC}\alpha^{(i)}\pi\delta) * \cos(2J_{NC}\alpha^{(i-1)}\pi\delta) * \cos(1J_{NH}\pi\delta) * \cos(\Omega_N t_1) \end{aligned}$$

$$\hat{H} = \text{echo (hetero)} \quad \hat{H} = C_z(\Omega_{C\alpha}[t_2]) \quad \downarrow$$

at the end of t₂

$$\begin{aligned} -4\mathbf{H}_y\mathbf{N}_x\mathbf{C}\alpha^{(i-1)}_y & * \cos(1J_{NC}\alpha^{(i)}\pi\delta) * \sin(2J_{NC}\alpha^{(i-1)}\pi\delta) * \cos(1J_{NH}\pi\delta) * \cos(\Omega_N t_1) * \cos(\Omega_C \alpha^{(i-1)} t_2) \\ -4\mathbf{H}_y\mathbf{N}_x\mathbf{C}\alpha^{(i)}_y & * \sin(1J_{NC}\alpha^{(i)}\pi\delta) * \cos(2J_{NC}\alpha^{(i-1)}\pi\delta) * \cos(1J_{NH}\pi\delta) * \cos(\Omega_N t_1) * \cos(\Omega_C \alpha^{(i)} t_2) \end{aligned}$$

memo: Both of the two 3-spin terms contain the $\sin(\Omega_C \alpha^{(i-1)} t_2)$ modulation (not shown above) were phase cycled out ($\varphi 2$).

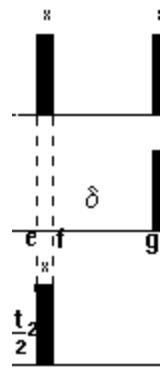
memo: 180° applied on both ¹H and ¹⁵N refocuses their chemical shifts (echo), and thus the evolution of $\mathbf{H}_y\mathbf{N}_x\mathbf{C}\alpha_y$ term depends only on the chemical shift of C_α, therefore the effective $\hat{H} = C_z(\Omega_{C\alpha}[t_2])$.

Note however, that because the $1J_{C\alpha C\beta}$ coupling is active during t₂, the $\cos(1J_{C\alpha C\beta}\pi t_2)$ modulation is effective, at the end of t₂. Thus, at $\sigma[e]$ is :

$$\begin{aligned} -4\mathbf{H}_y\mathbf{N}_x\mathbf{C}\alpha^{(i-1)}_y & * \cos(1J_{NC}\alpha^{(i)}\pi\delta) * \sin(2J_{NC}\alpha^{(i-1)}\pi\delta) * \cos(1J_{NH}\pi\delta) * \cos(\Omega_N t_1) * \cos(\Omega_C \alpha^{(i-1)} t_2) \\ & * \cos(1J_C \alpha_C \beta \pi t_2) \\ -4\mathbf{H}_y\mathbf{N}_x\mathbf{C}\alpha^{(i)}_y & * \sin(1J_{NC}\alpha^{(i)}\pi\delta) * \cos(2J_{NC}\alpha^{(i-1)}\pi\delta) * \cos(1J_{NH}\pi\delta) * \cos(\Omega_N t_1) * \cos(\Omega_C \alpha^{(i)} t_2) \\ & * \cos(1J_C \alpha_C \beta \pi t_2) \end{aligned}$$

F. module (back transforming the multiple-quantum coherences to an anti-phase coherence on N)

$$\begin{array}{l}
 \sigma[e] \\
 \hat{H} = H_x(\pi/2) \\
 \\
 \hat{H} = C\alpha_x(\pi/2) \\
 \end{array}
 \begin{array}{c}
 -4H_yN_xC_y \\
 \downarrow \\
 +4H_zN_xC_y \\
 \downarrow \\
 -4H_zN_xC_z
 \end{array}$$



σ[f]

$$\begin{aligned}
 & -4H_zN_xC\alpha^{(i-1)}_z * \cos(^1J_{NC}\alpha^{(i)}\pi\delta) * \sin(^2J_{NC}\alpha^{(i-1)}\pi\delta) * \cos(^1J_{NH}\pi\delta) * \cos(\Omega_N t_1) * \cos(\Omega_C\alpha^{(i-1)}t_2) \\
 & * \cos(^1J_C\alpha_C\beta\pi t_2) \\
 & -4H_zN_xC\alpha^{(i)}_z * \sin(^1J_{NC}\alpha^{(i)}\pi\delta) * \cos(^2J_{NC}\alpha^{(i-1)}\pi\delta) * \cos(^1J_{NH}\pi\delta) * \cos(\Omega_N t_1) * \cos(\Omega_C\alpha^{(i)}t_2) \\
 & * \cos(^1J_C\alpha_C\beta\pi t_2)
 \end{aligned}$$

G. module (back transferring coherence from Cα(s) to N during the second delay-time δ)

$$\begin{array}{l}
 \sigma[f] \\
 \hat{H} = 2N_zC_z(^1J_{NC}\alpha^{(i)}\pi\delta) \\
 \\
 \hat{H} = 2N_zC_z(^2J_{NC}\alpha^{(i-1)}\pi\delta) \\
 -4H_zN_xC\alpha^{(i)}_z
 \end{array}
 \begin{array}{c}
 -4H_zN_xC_z \\
 \downarrow \\
 -4H_zN_xC\alpha^{(i)}_z \\
 \downarrow \\
 -8H_zN_yC\alpha^{(i)}_zC\alpha^{(i-1)}_z
 \end{array}
 \begin{array}{c}
 -2H_zN_y \\
 \downarrow \\
 -2H_zN_y
 \end{array}
 \begin{array}{c}
 +4H_zN_xC\alpha^{(i-1)}_z
 \end{array}$$

memo : The reverse INEPT pulses (as it is the terminating module of this pulse sequence) has to result in observable magnetization, thus among the 4 terms:

- $4\mathbf{H}_z \mathbf{N}_x \mathbf{C}\alpha^{(i)}_z \rightarrow -4\mathbf{H}_y \mathbf{N}_x \mathbf{C}\alpha^{(i)}_z$ is a multiple-quantum coh. → **not** observable
- + $4\mathbf{H}_z \mathbf{N}_x \mathbf{C}\alpha^{(i-1)}_z \rightarrow +4\mathbf{H}_y \mathbf{N}_x \mathbf{C}\alpha^{(i-1)}_z$ is a multiple-quantum coh. → **not** observable
- $8\mathbf{H}_z \mathbf{N}_y \mathbf{C}\alpha^{(i)}_z \mathbf{C}\alpha^{(i-1)}_z \rightarrow -4\mathbf{H}_x \mathbf{C}\alpha^{(i)}_z \mathbf{C}\alpha^{(i-1)}_z$ is an anti-phase coh. → observable **but ignored**, as it is modulated by unwanted proton-carbon couplings: ${}^2J_{HC}\alpha^{(i)}$ and ${}^3J_{HC}\alpha^{(i)}$
- $2\mathbf{H}_z \mathbf{N}_y \rightarrow -\mathbf{H}_x$ the only important term to be considered**

Therefore from:

$\sigma[f]$

$$\begin{aligned}
 & -4\mathbf{H}_z \mathbf{N}_x \mathbf{C}\alpha^{(i-1)}_z * \cos({}^1J_{NC}\alpha^{(i)} \pi\delta) * \sin({}^2J_{NC}\alpha^{(i-1)} \pi\delta) * \cos({}^1J_{NH} \pi\delta) * \cos(\Omega_N t_1) * \cos(\Omega_C \alpha^{(i-1)} t_2) \\
 & * \cos({}^1J_C \alpha_C \beta \pi t_2) \\
 & -4\mathbf{H}_z \mathbf{N}_x \mathbf{C}\alpha^{(i)}_z * \sin({}^1J_{NC}\alpha^{(i)} \pi\delta) * \cos({}^2J_{NC}\alpha^{(i-1)} \pi\delta) * \cos({}^1J_{NH} \pi\delta) * \cos(\Omega_N t_1) * \cos(\Omega_C \alpha^{(i)} t_2) \\
 & * \cos({}^1J_C \alpha_C \beta \pi t_2)
 \end{aligned}$$

one gets the following two terms:

$\sigma[g]$

$$\begin{aligned}
 & -2\mathbf{H}_z \mathbf{N}_y * \boxed{\sin({}^1J_{NC}\alpha^{(i)} \pi\delta) * \cos({}^2J_{NC}\alpha^{(i-1)} \pi\delta)} * \cos({}^1J_{NC}\alpha^{(i)} \pi\delta) * \sin({}^2J_{NC}\alpha^{(i-1)} \pi\delta) * \cos({}^1J_{NH} \pi\delta) \\
 & * \cos(\Omega_N t_1) * \cos(\Omega_C \alpha^{(i-1)} t_2) * \cos({}^1J_C \alpha_C \beta \pi t_2) \\
 & -2\mathbf{H}_z \mathbf{N}_y * \boxed{\sin({}^1J_{NC}\alpha^{(i)} \pi\delta) * \cos({}^2J_{NC}\alpha^{(i-1)} \pi\delta)} * \sin({}^1J_{NC}\alpha^{(i)} \pi\delta) * \cos({}^2J_{NC}\alpha^{(i-1)} \pi\delta) * \cos({}^1J_{NH} \pi\delta) \\
 & * \cos(\Omega_N t_1) * \cos(\Omega_C \alpha^{(i)} t_2) * \cos({}^1J_C \alpha_C \beta \pi t_2)
 \end{aligned}$$

$\sigma[g]$

$$\begin{aligned}
 & -2\mathbf{H}_z \mathbf{N}_y * \sin(1J_{NC}\alpha^{(i)}\pi\delta) * \cos(2J_{NC}\alpha^{(i-1)}\pi\delta) * \cos(1J_{NC}\alpha^{(i)}\pi\delta) * \sin(2J_{NC}\alpha^{(i-1)}\pi\delta) * \cos^2(1J_{NH}\pi\delta) \\
 & * \cos(\Omega_N t_1) * \cos(\Omega_C \alpha^{(i-1)} t_2) * \cos(1J_C \alpha_C \beta \pi t_2) \\
 & -2\mathbf{H}_z \mathbf{N}_y * \sin(1J_{NC}\alpha^{(i)}\pi\delta) * \cos(2J_{NC}\alpha^{(i-1)}\pi\delta) * \sin(1J_{NC}\alpha^{(i)}\pi\delta) * \cos(2J_{NC}\alpha^{(i-1)}\pi\delta) * \cos^2(1J_{NH}\pi\delta) \\
 & * \cos(\Omega_N t_1) * \cos(\Omega_C \alpha^{(i)} t_2) * \cos(1J_C \alpha_C \beta \pi t_2)
 \end{aligned}$$

H. module (finishing with a reverse INEPT)

$\sigma[g]$ "The reverse INEPT module"

$$\hat{H} = H_x(\pi/2)$$

$$\hat{H} = N_x(\pi/2)$$

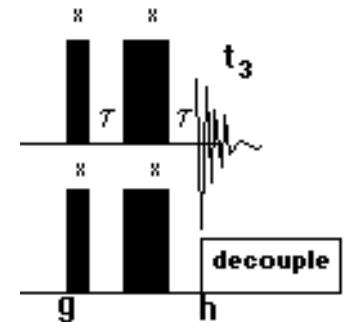
$$\hat{H} = \text{echo (homo)}$$

$$\{-H_x \sin(J_{HN}\pi 2\tau) + 2H_y N_z \cos(J_{HN}\pi 2\tau) \text{ with } 2\tau = 1/2J_{HN}\}$$

$\sigma[h]$ „at the end, before acquisition (t_3) starts,,: $-H_x$

Note that this final observable in-phased $-H_x$ coherence (or transverse magnetization) is modulated as follows:

$$\begin{aligned}
 & -H_x * \sin(1J_{NC}\alpha^{(i)}\pi\delta) * \cos(2J_{NC}\alpha^{(i-1)}\pi\delta) * \cos(1J_{NC}\alpha^{(i)}\pi\delta) * \sin(2J_{NC}\alpha^{(i-1)}\pi\delta) * \cos^2(1J_{NH}\pi\delta) \\
 & * \cos(\Omega_N t_1) * \cos(\Omega_C \alpha^{(i-1)} t_2) * \cos(1J_C \alpha_C \beta \pi t_2) \\
 & -H_x * \sin(1J_{NC}\alpha^{(i)}\pi\delta) * \cos(2J_{NC}\alpha^{(i-1)}\pi\delta) * \sin(1J_{NC}\alpha^{(i)}\pi\delta) * \cos(2J_{NC}\alpha^{(i-1)}\pi\delta) * \cos^2(1J_{NH}\pi\delta) \\
 & * \cos(\Omega_N t_1) * \cos(\Omega_C \alpha^{(i)} t_2) * \cos(1J_C \alpha_C \beta \pi t_2)
 \end{aligned}$$



ACQ:

As no coupling is effective during acquisition

$$\hat{H} = H_z(\Omega_H[t_3]) \quad \begin{matrix} -H_x \\ \downarrow \\ -H_x \cos(\Omega_H t_3) + -H_y \sin(\Omega_H t_3) \end{matrix}$$

memo : put the receiver on -x

therefore only the single x term remains:

the following term can be found

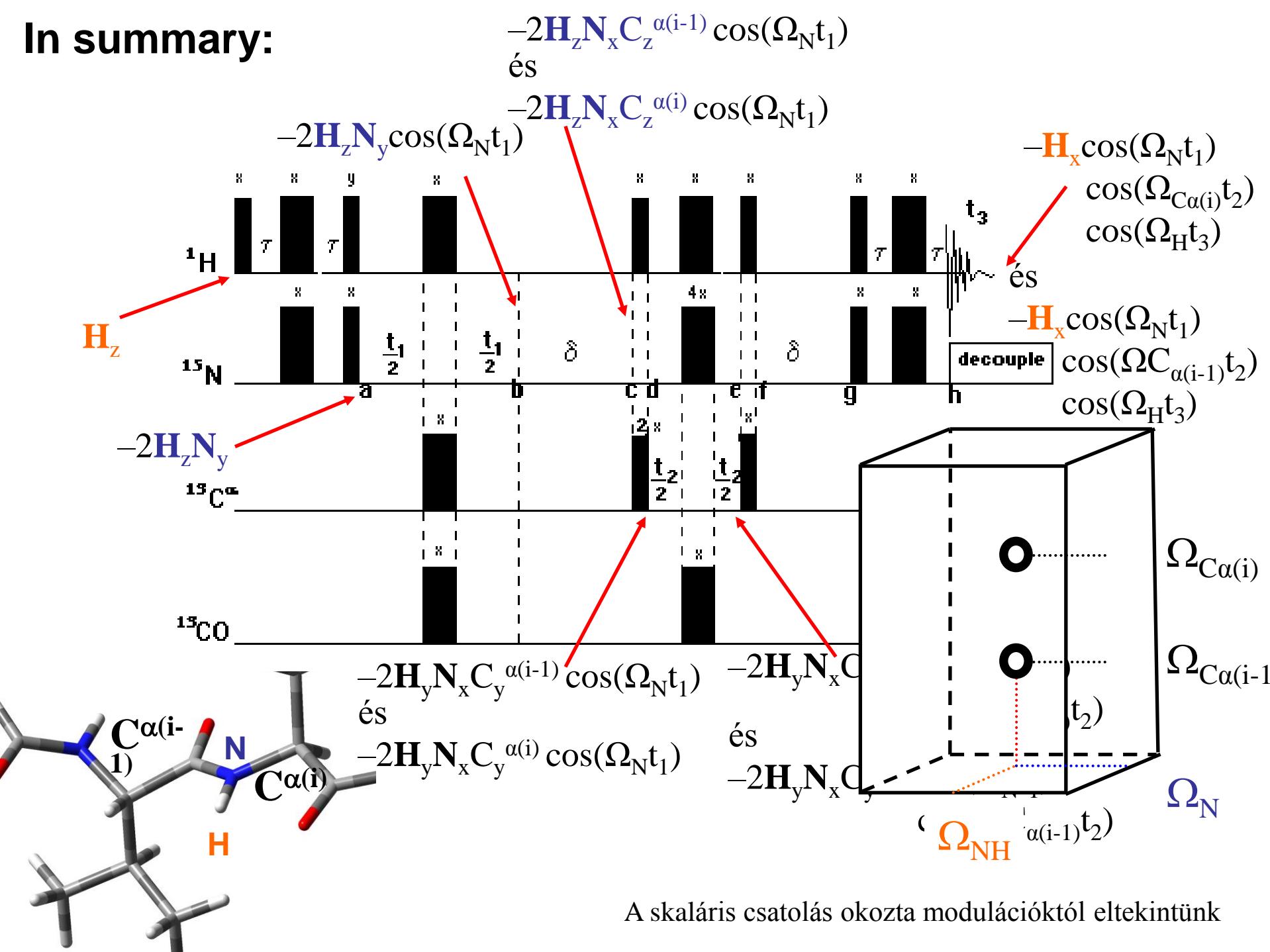
$$\begin{matrix} +H_x \cos(\Omega_H t_3) \\ H_x [+ \text{ at } \Omega_H, \Omega_N, \Omega_C \alpha] \end{matrix}$$

if one sets the phase that *cos* is absorptive in all three dimensions (t_3, t_2, t_1),
then:

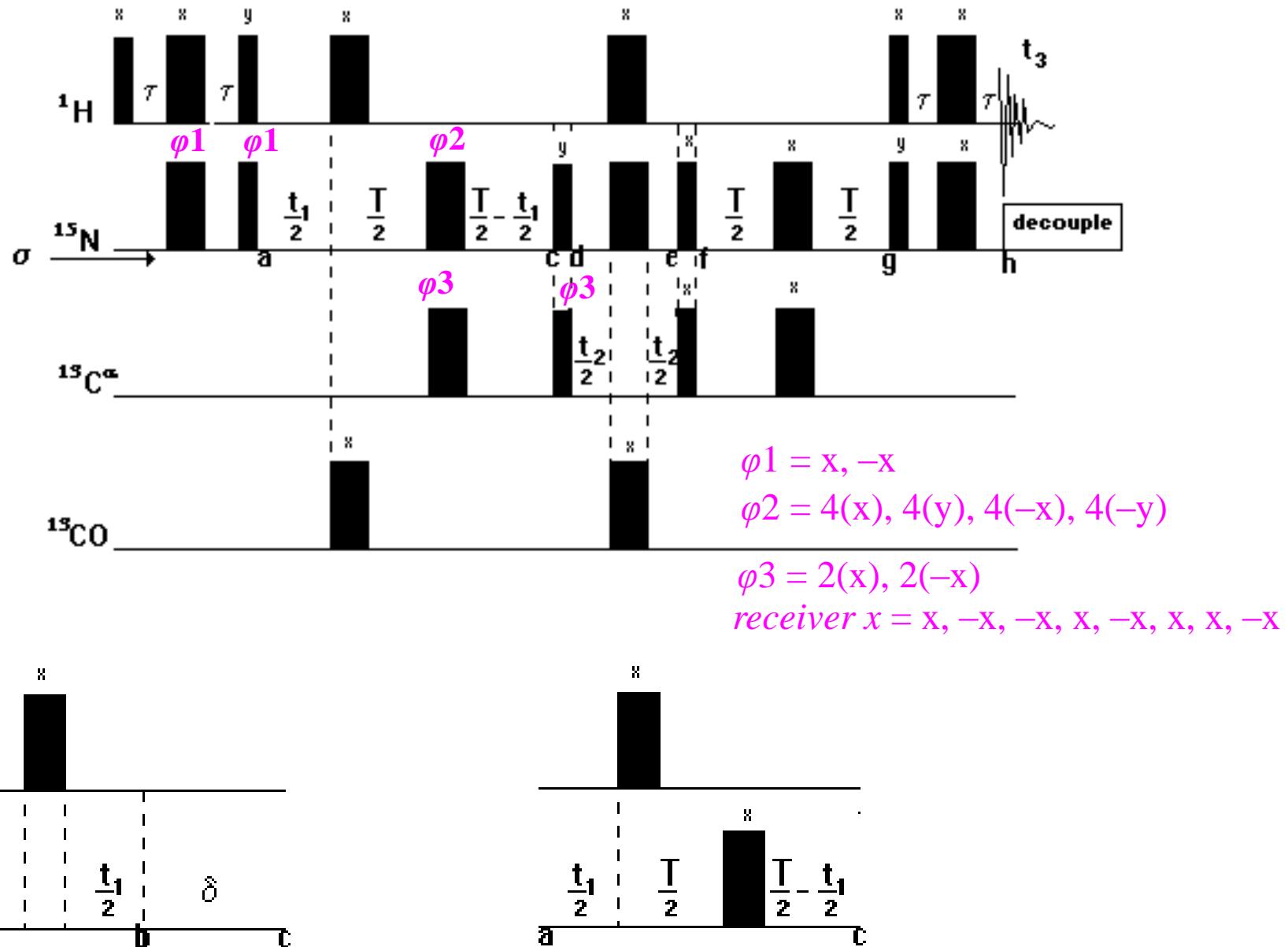
$$\begin{matrix} H_x [+a \text{ at } \Omega_H, \Omega_N, \Omega_C \alpha^{(i-1)}] \\ H_x [+a \text{ at } \Omega_H, \Omega_N, \Omega_C \alpha^{(i)}] \end{matrix}$$

$$\begin{aligned} & -H_x * \sin(1J_{NC} \alpha^{(i)} \pi \delta) * \cos(2J_{NC} \alpha^{(i-1)} \pi \delta) * \cos(1J_{NC} \alpha^{(i)} \pi \delta) * \sin(2J_{NC} \alpha^{(i-1)} \pi \delta) * \cos^2(1J_{NH} \pi \delta) \\ & * \cos(1J_C \alpha_C \beta \pi t_2) * \cos(\Omega_N t_1) * \cos(\Omega_C \alpha^{(i-1)} t_2) * \cos(\Omega_H t_3) \\ & -H_x * \sin(1J_{NC} \alpha^{(i)} \pi \delta) * \cos(2J_{NC} \alpha^{(i-1)} \pi \delta) * \sin(1J_{NC} \alpha^{(i)} \pi \delta) * \cos(2J_{NC} \alpha^{(i-1)} \pi \delta) * \cos^2(1J_{NH} \pi \delta) \\ & * \cos(1J_C \alpha_C \beta \pi t_2) * \cos(\Omega_N t_1) * \cos(\Omega_C \alpha^{(i)} t_2) * \cos(\Omega_H t_3) \end{aligned}$$

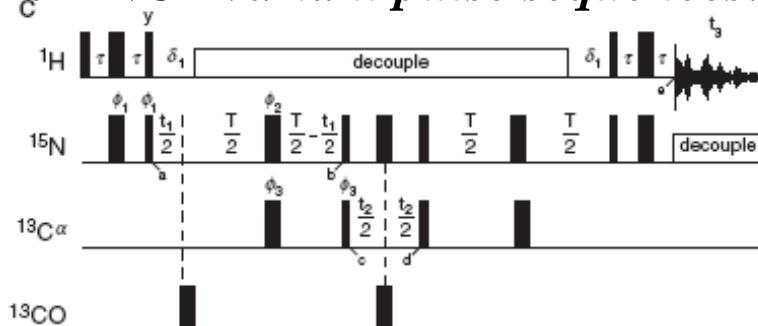
In summary:



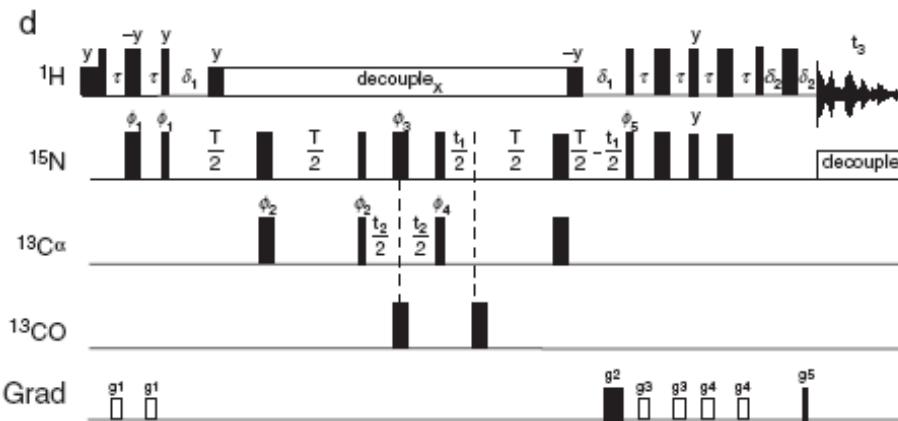
A constant-time HNCA (CT-HNCA) variant pulse sequence:



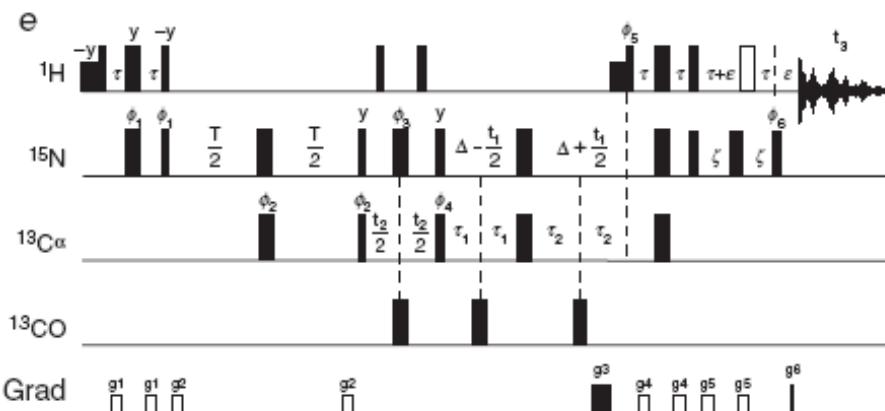
More HNCA variant pulse sequences:



Decoupled-CT-HNCA
synchronous broad band decoupling
(e.g. WALTZ-16) enhances the
sensitivity of a CT-HNCA.



PFG-PEP-HNCA
at the end the reversed-INEPT
module is replaced by a PEP-
reversed-INEPT.



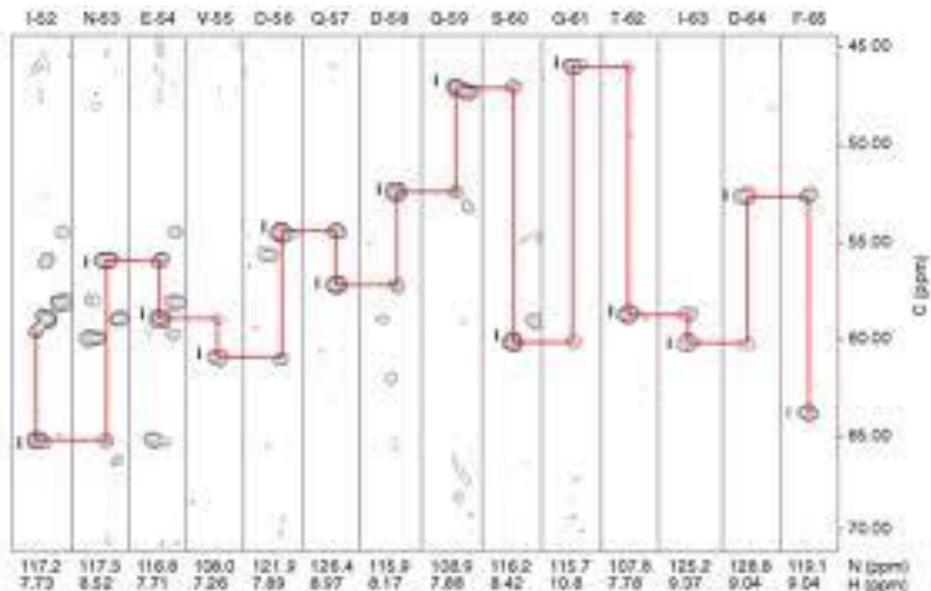
PFG-TROSY-HNCA

evolution of ^{15}N and
the coherence transfer from N to Ca(s)
are **separated** in time

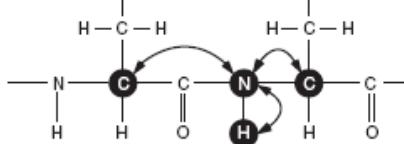
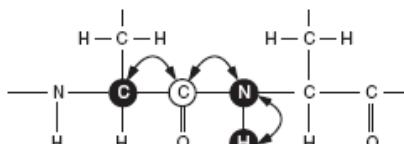
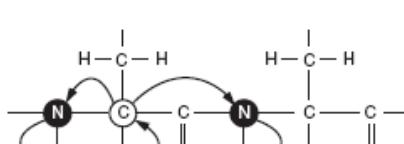
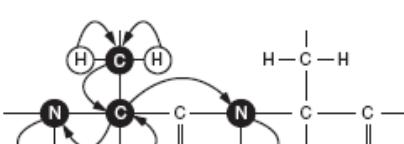
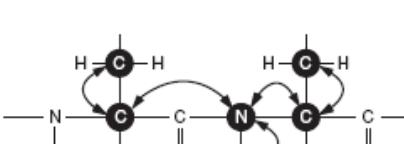
evolution of ^{15}N and
the coherence transfer from N to Ca(s)
are **simultaneous**

advantage: one saves time with the CT- experiment.

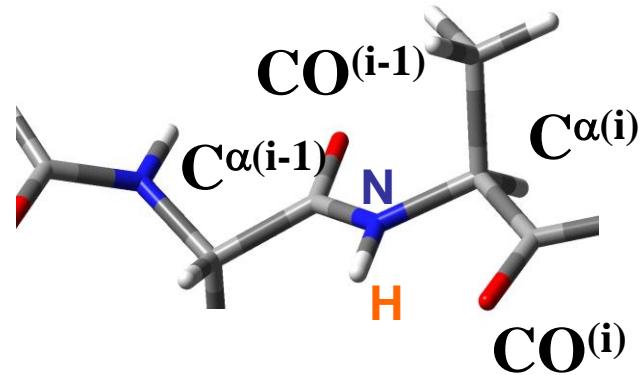
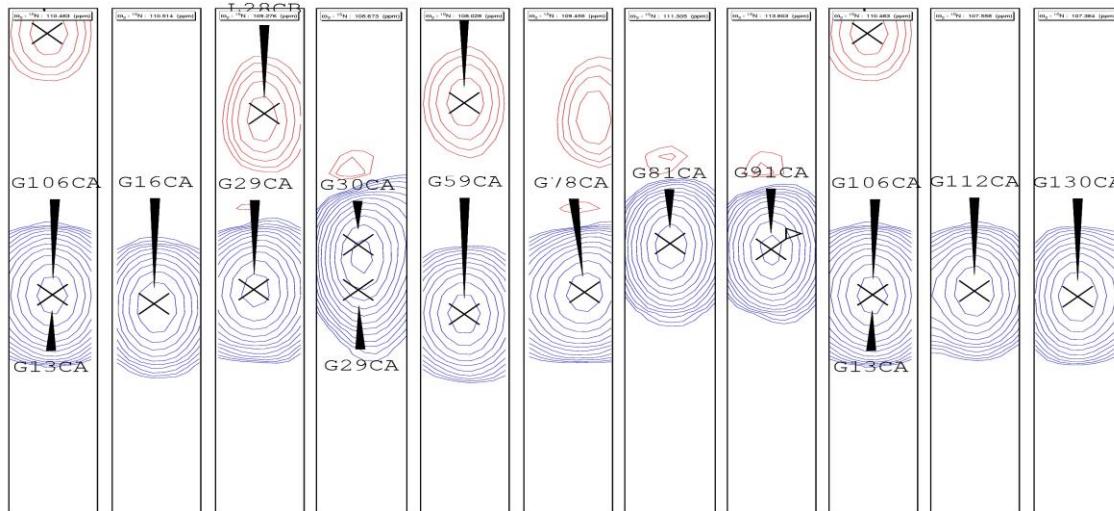
memo: by saving time one may gain in sensitivity as relaxation is reduced.



The simplest triple resonance experiments used for backbone resonance assignment

Experiment	Correlations observed	Magnetization transfer	J couplings ^b
HNCA	$^1\text{H}_i^{\text{N}} - ^{15}\text{N}_i - ^{13}\text{C}_i^{\alpha}$ $^1\text{H}_i^{\text{N}} - ^{15}\text{N}_i - ^{13}\text{C}_{i-1}^{\alpha}$		$^1\text{J}_{\text{NH}}$ $^1\text{J}_{\text{NC}^{\alpha}}$ $^2\text{J}_{\text{NC}^{\alpha}}$
HN(CO)CA	$^1\text{H}_i^{\text{N}} - ^{15}\text{N}_i - ^{13}\text{C}_{i-1}^{\alpha}$		$^1\text{J}_{\text{NH}}$ $^1\text{J}_{\text{NCO}}$ $^1\text{J}_{\text{C}^{\alpha}\text{CO}}$
H(CA)NH	$^1\text{H}_i^{\alpha} - ^{15}\text{N}_i - ^1\text{H}_i^{\text{N}}$ $^1\text{H}_i^{\alpha} - ^{15}\text{N}_{i+1} - ^1\text{H}_{i+1}^{\text{N}}$		$^1\text{J}_{\text{C}^{\alpha}\text{H}^{\alpha}}$ $^1\text{J}_{\text{NC}^{\alpha}}$ $^2\text{J}_{\text{NC}^{\alpha}}$ $^1\text{J}_{\text{NH}}$
CBCANH	$^{13}\text{C}_i^{\beta}/^{13}\text{C}_i^{\alpha} - ^{15}\text{N}_i - ^1\text{H}_i^{\text{N}}$ $^{13}\text{C}_i^{\beta}/^{13}\text{C}_i^{\alpha} - ^{15}\text{N}_{i+1} - ^1\text{H}_{i+1}^{\text{N}}$		$^1\text{J}_{\text{CH}}$ $^1\text{J}_{\text{C}^{\alpha}\text{C}^{\beta}}$ $^1\text{J}_{\text{NC}^{\alpha}}$ $^2\text{J}_{\text{NC}^{\alpha}}$ $^1\text{J}_{\text{NH}}$
HNCACB	$^{13}\text{C}_i^{\beta}/^{13}\text{C}_i^{\alpha} - ^{15}\text{N}_i - ^1\text{H}_i^{\text{N}}$ $^{13}\text{C}_{i-1}^{\beta}/^{13}\text{C}_{i-1}^{\alpha} - ^{15}\text{N}_i - ^1\text{H}_i^{\text{N}}$		$^1\text{J}_{\text{C}^{\alpha}\text{C}^{\beta}}$ $^1\text{J}_{\text{NC}^{\alpha}}$ $^2\text{J}_{\text{NC}^{\alpha}}$ $^1\text{J}_{\text{NH}}$

Resonance assignment in ^{15}N - ^{13}C -Calpastatin



HNCA and HN(CO)CA

